w-singularities in cosmological models

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Spanish Relativity Meeting 2010



Sketch

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Motivation

- Astronomical observations of luminosity distances derived from Type Ia supernovae, CMB spectrum and global matter distribution provide evidence of cosmic speed up of the Universe.
- One approach to deal with this riddle is to consider modifications of general relativity.
- Alternatively, cosmic acceleration might be due to an exotic fluid filling the Universe, known as dark energy.
- These have given rise to a collection of new cosmological evolutions, future singularites being the most perplexing ones ("big rip", "sudden singularities"...).

Classification of singularities

- Big Bang / Crunch: zero a, divergent H, density and pressure. Strong.
- Type I: "Big Rip": divergent a. Strong.
- Type II: "Sudden": finite a, H, density, divergent H and pressure. Weak
- Type III: "Big Freeze": finite a, divergent H, density and pressure. Weak/Strong
- Type IV: "Big Brake": finite *a*, *H*, *H*, density and pressure but divergent higher derivatives. Weak.



w-singularities

Dabrowski & Denkiewicz introduced a cosmological model,

$$a(t) = \frac{a_s}{1 - \frac{3\gamma}{2} \left(\frac{n-1}{n - \frac{2}{3\gamma}}\right)^{n-1}} + \frac{1 - \frac{2}{3\gamma}}{n - \frac{2}{3\gamma}} \frac{na_s}{1 - \frac{2}{3\gamma} \left(\frac{n - \frac{2}{3\gamma}}{n-1}\right)^{n-1}} \left(\frac{t}{t_s}\right)^{\frac{2}{3\gamma}}$$

$$+ \frac{a_s}{\frac{3\gamma}{2} \left(\frac{n-1}{n - \frac{2}{3\gamma}}\right)^{n-1}} \left(1 - \frac{1 - \frac{2}{3\gamma}}{n - \frac{2}{3\gamma}} \frac{t}{t_s}\right)^n, \quad \gamma > 0, \quad n \neq 1.$$

- The scale factor $a(t_s) = a_s$ is finite.
- Density and pressure vanish at t_s , $\rho(t_s) = 0 = \rho(t_s)$.
- If *n* is natural, the derivatives of *a* and the Hubble parameter are regular.
- Just the barotropic index w ($p = w\rho$) is singular at $t = t_s$.



Barotropic index expansions

The barotropic index *w* may be written in terms of the scale factor,

$$w=-\frac{1}{3}-\frac{2}{3}\frac{\ddot{a}\ddot{a}}{\dot{a}^2}.$$

Assuming that the scale factor admits a generalized power expansion with real exponents,

$$a(t) = c_0 (t_s - t)^{\eta_0} + c_1 (t_s - t)^{\eta_1} + \cdots, \quad \eta_0 < \eta_1 < \cdots, \quad c_0 > 0,$$

η_0	η_1	η_2	W _S
<i>≠</i> 0	(η_0,∞)	(η_1,∞)	Finite
0	(0,1)	(η_1,∞)	Infinite
	1	(1,2)	Infinite
	1	$[2,\infty)$	Finite
	$(1,\infty)$	$[\eta_1,\infty)$	Infinite



Density and pressure

Besides a diverging barotropic index, we need vanishing density and pressure.

For
$$\eta_0 = 0$$
, $\eta_1 \neq 1$,

$$\rho = \frac{3c_1^2\eta_1^2}{c_0^2}(t_s-t)^{2(\eta_1-1)}+\cdots,
\rho = -\frac{2c_1\eta_1(\eta_1-1)}{c_0}(t_s-t)^{\eta_1-2}+\cdots,$$

 $\eta_1 > 2$ is required.



Results

■ A FLRW cosmological model has a singular barotropic index w with vanishing pressure and density at t_s if an only if the generalized power expansion of a(t) is of the form

$$a(t) = c_0 + c_1(t_s - t)^{\eta_1} + \cdots,$$

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■ A FLRW cosmological model has a w-singularity at a finite time t_s if an only if the scale factor a(t) admits a Taylor series at t_s with vanishing linear and quadratic terms,

$$a(t) = c_0 + O((t - t_s)^3).$$



Scale factor vs. singularities

η_0	η_1	η_2	Tipler	Królak	N.O.T.
(-s,0)	(η_0,s)	(η_1,s)	Strong	Strong	1
0	(0, 1)	(η_1,s)	Weak	Strong	III
	(1, 2)	(η_1,s)	Weak	Weak	II
	1	(1, 2)	Weak	Weak	II
		[2, s)	Weak	Weak	IV
	[2, s)	(η_1,s)	Weak	Weak	IV
(0, s)	(η_0,s)	(η_1,s)	Strong	Strong	Crunch

Hence w-singularities are weak singularities.



Further reading

- L. Fernández-Jambrina,
 On w-cosmological singularities (preprint)
- L. Fernández-Jambrina, R. Lazkoz, Singular fate of the universe in modified theories of gravity Phys. Lett. B 670 254 (2009) [arXiv:0805.2284]
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 Phys. Rev. D 74 064030 (2006) [arXiv:gr-qc/0607073]
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Thank you all!



