

Principles of Einstein–Finsler Gravity and Perspectives in Modern Cosmology

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Goals

- **Finsler modifications of GR** derived for QG theories; Geometric models for quantum contributions and **LV**
- Canonical models for **Einstein–Finsler gravity** (EFG); principles and axioms
- Physical implications of EFG: **Finsler branes**, locally anisotropic cosmology & astrophysics

Reviews and new results:

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Motivation: nonlinear disp; QG & LV, cosmology

1. Deforms in Minkovski s-t: $E^2 = p^2 c^2 + m_0^2 c^4 + \varphi(E, p; \mu; M_P)$

$$E \sim \frac{\partial \phi}{\partial t}, p_{\hat{i}} \sim \frac{\partial \phi}{\partial x^{\hat{i}}}, \omega = \frac{\partial \phi}{\partial t} \quad k_{\hat{i}} = \frac{\partial \phi}{\partial x^{\hat{i}}}, \omega^2 = c_s^2 k^2 + c_s^2 \left(\frac{\hbar}{2m_0 c_s} \right)^2 k^4 + \dots$$

effective c_s , ($x^1 = ct, x^2, x^3, x^4$); $\hat{i}, \hat{j} \dots = 2, 3, 4$;

$$\omega^2 = c^2 [g_{\hat{i}\hat{j}} k^{\hat{i}} k^{\hat{j}}]^2 (1 - q_{\hat{i}_1 \hat{i}_2 \dots \hat{i}_{2r}} y^{\hat{i}_1} \dots y^{\hat{i}_{2r}} / r [g_{\hat{i}\hat{j}} k^{\hat{i}} k^{\hat{j}}]^{2r})$$

light velocity in "media/ether" $c^2 = g_{\hat{i}\hat{j}}(x^i) y^{\hat{i}} y^{\hat{j}} / \tau^2 \rightarrow \check{F}^2(y^{\hat{i}}) / \tau^2$

fundamental Finsler function $F(x^i, \beta y^j) = \beta F(x^i, y^j), \beta > 0$,

$$ds^2 = F^2 \approx -(cdt)^2 + g_{\hat{i}\hat{j}}(x^k) y^{\hat{i}} y^{\hat{j}} \left[1 + \frac{1}{r} \frac{q_{\hat{i}_1 \hat{i}_2 \dots \hat{i}_{2r}}(x^k) y^{\hat{i}_1} \dots y^{\hat{i}_{2r}}}{(g_{\hat{i}\hat{j}}(x^k) y^{\hat{i}} y^{\hat{j}})^r} \right] + O(q^2)$$

Finsler "metrics", velocities on TV , $F g_{ij}(x^i, y^j) = \frac{1}{2} \frac{\partial F^2}{\partial y^i \partial y^j}$

2. Exact solutions & modified cosmology with generic off-diagonal metrics and local anisotropy.

Einstein–Finsler Gravity (EFG)

Statement I: A (pseudo) Finsler metric, ${}^F g_{ij}(x^k, y^a)$, **DOES NOT** define completely a geometric model (not Riemannian !)

Statement II: A model of Finsler geometry is defined on TV by **THREE** fundamental geometric objects induced by $F(x, y)$:

- ① **N–connection**, $N_i^a(x, y)$, splitting ${}^F \mathbf{N} : TTV = hTV \oplus vTV$ canonically, Euler–Lagrange for $L = F^2$ are semi–sprays,
- ② **d–connection**, N–adapted linear connect. ${}^F \mathbf{D} = (hD, vD)$, preferred/ canonically induced by ${}^F g_{ij}$ and N_i^a
- ③ **d–metric**, ${}^F \mathbf{g} = hg \oplus vg$

2 classes: a) nonmetricity, ${}^F \mathbf{Q} := {}^F \mathbf{D} {}^F \mathbf{g}$, **Chern d–conn.**, ${}^{Ch} \mathbf{D}$

b) metricity, ${}^F \mathbf{Q} = 0$, **Cartan d–conn.**, ${}^{Cart} \mathbf{D}$

Levi–Civita ${}^F \nabla$ is **NOT** adapted to nonholonomic ${}^F \mathbf{N}$.

\exists induced by ${}^F \mathbf{g}$: torsion ${}^F \mathbf{T}$, and/or ${}^F \mathbf{Q}$ (not Riemann–Cartan)



Einstein–Finsler spacetimes/gravity, EFG

Spacetime as a **nonholonomic manifold**/ bundle $\mathbf{V} := (V, \mathcal{D})$ (Vrănceanu, 1926), or TM , with a non–integrable distribution \mathcal{D} .

Geometric data: Finsler $(F : \mathbf{N}, \mathbf{D}, \mathbf{g})$ and Riemannian (∇, \mathbf{g})

N–anholonomic frames: $\mathbf{e}_\nu = (\mathbf{e}_i = \partial_i - N_i^a \partial_a, \mathbf{e}_a = \partial_a)$

Sasaki d–metric: ${}^F \mathbf{g} = {}^F g_{ij}(u) dx^i \otimes dx^j + {}^F g_{ab}(u) {}^c \mathbf{e}^a \otimes {}^c \mathbf{e}^b$,
 for ${}^c \mathbf{e}^a = dy^a + {}^c N_i^a(u) dx^i$.

For $\tilde{\mathbf{D}}$, standard Riemannian, Ricci, Einstein d–tensors; h–/v–splitting.

N–adapted coef.: ${}^{Cart} \mathbf{D} = \tilde{\mathbf{D}} = (h\tilde{D}, v\tilde{D}) = \{\tilde{\Gamma}_{\gamma\tau}^\alpha = (\tilde{L}_{jk}^i, \tilde{C}_{bc}^a)\}$,

$$\tilde{L}_{jk}^i = \frac{1}{2} {}^F g^{ir} (\mathbf{e}_k {}^F g_{jr} + \mathbf{e}_j {}^F g_{kr} - \mathbf{e}_r {}^F g_{jk}),$$

$$\tilde{C}_{bc}^a = \frac{1}{2} {}^F g^{ad} (e_c {}^F g_{bd} + e_c {}^F g_{cd} - e_d {}^F g_{bc}).$$

Theorem: Equivalent (pseudo) Finsler & Riemannian theories
 if ${}^g \mathbf{D} = {}^g \nabla + {}^g \mathbf{Z}$, distortion determined by $\mathbf{g} = {}^F \mathbf{g}$.

Principles and axioms of EFG

Principles: Similarly to GR with ${}^g\nabla$ on V construct **EFG:** with $\mathbf{g} \sim {}^F\mathbf{g}$, $\mathbf{N} \sim {}^F\mathbf{N}$ and ${}^{Cart}\mathbf{D}$ on TV , or \mathbf{V} .

- 1 **Generalized equivalence principle:** Ideas on Free Fall and Universality of Gravitational Redshift for ${}^{Cart}\mathbf{D}$.
- 2 **Generalized Mach principle:** quantum energy/motion encoded via $(\mathbf{N}, \mathbf{g}, \mathbf{D})$ for spacetime ether with y^a .
- 3 **Principle of general covariance** extended on \mathbf{V} , or TV , with "mixing of Finsler parametrizations".
- 4 **Motion eqs and conservation laws:** Nonholonomic Bianchi identities for ${}^F\mathbf{D}$; $\nabla_i T^{ij} = 0 \rightarrow \mathbf{D}_\alpha \Upsilon^{\alpha\beta} \neq 0$.
- 5 **Einstein–Finsler gravitational field eqs** for ${}^F\mathbf{D}$.
- 6 **Axiomatics:** Constructive–axiomatic appr. (Ehlers–Pirani–Schild, EPS axioms), paradigm "Lorentzian 4–manifold" in GR; nonholon. tangent bundle on "L ..." for EFG.

Gravitational field eqs in EFG

$\forall \mathbf{D}$, Einstein eqs: $\mathbf{E}_{\alpha\beta} = \Upsilon_{\alpha\beta}$,

h -/ v -components, for $R_{ai} = R^b_{aib}$ and $R_{ia} = R^k_{ikb}$:

$$R_{ij} - \frac{1}{2}(R + S)g_{ij} = \Upsilon_{ij},$$

$$R_{ab} - \frac{1}{2}(R + S)h_{ab} = \Upsilon_{ab},$$

$$R_{ai} = \Upsilon_{ai}, \quad R_{ia} = -\Upsilon_{ia},$$

Remark: For $^{Cart}\mathbf{D}$, general off-diagonal solutions for EFG,

restrictions to GR, $\mathbf{g} = \underline{g}_{\alpha\beta}(u) du^\alpha \otimes du^\beta$,

$$\underline{g}_{\alpha\beta} = \begin{bmatrix} g_{ij} + N_i^a N_j^b h_{ab} & N_j^e h_{ae} \\ N_i^e h_{be} & h_{ab} \end{bmatrix}, \text{ where } N_i^a \neq A_{bi}^a(x) y^b$$

Claim: Compactification/trapping/warping mechanism on velocity/momenta for a "new" QG and LV phenomenology.

Finsler–branes & cosmological solutions

Nonholon. trapping solutions (cosmology, with $h_3(x^i, y^3 = t)$) :

$$\mathbf{g} = g_1 dx^1 \otimes dx^1 + g_2 dx^2 \otimes dx^2 + h_3 \mathbf{e}^3 \otimes \mathbf{e}^3 + h_4 \mathbf{e}^4 \otimes \mathbf{e}^4 + (l_P)^2 \frac{\bar{h}}{\phi^2} [{}^q h_5 \mathbf{e}^5 \otimes \mathbf{e}^5 + {}^q h_6 \mathbf{e}^6 \otimes \mathbf{e}^6 + {}^q h_7 \mathbf{e}^7 \otimes \mathbf{e}^7 + {}^q h_8 \mathbf{e}^8 \otimes \mathbf{e}^8]$$

$$\mathbf{e}^3 = dy^3 + w_i dx^i, \mathbf{e}^4 = dy^4 + n_i dx^i, \mathbf{e}^5 = dy^5 + {}^1 w_i dx^i, \mathbf{e}^6 = dy^6 + {}^1 n_i dx^i, \mathbf{e}^7 = dy^7 + {}^2 w_i dx^i, \mathbf{e}^8 = dy^8 + {}^2 n_i dx^i.$$

$$\phi^2(y^5) = \frac{3\epsilon^2 + a(y^5)^2}{3\epsilon^2 + (y^5)^2} \text{ and } l_P \sqrt{|\bar{h}(y^5)|} = \frac{9\epsilon^4}{[3\epsilon^2 + (y^5)^2]^2},$$

N–connection coefficients determined by sources

$$\begin{aligned} {}^h \Lambda(x^i) &= \tilde{\Upsilon}_4 + \tilde{\Upsilon}_6 + \tilde{\Upsilon}_8, & {}^v \Lambda(x^i, v) &= \tilde{\Upsilon}_2 + \tilde{\Upsilon}_6 + \tilde{\Upsilon}_8, \\ {}^5 \Lambda(x^i, y^5) &= \tilde{\Upsilon}_2 + \tilde{\Upsilon}_4 + \tilde{\Upsilon}_8, & {}^7 \Lambda(x^i, y^5, y^7) &= \tilde{\Upsilon}_2 + \tilde{\Upsilon}_4 + \tilde{\Upsilon}_6. \end{aligned}$$

Conclusions

- **Almost all** models of QG with nonlinear dispersions can be geometrized as certain Finsler spacetimes.
- **Natural/ Canonical Principles** for metric compatible EFG generalizing the GR on $TV, \nabla \rightarrow^{Cart} \mathbf{D}$.
- Finsler branes, trapping: "new" QG/ LV phenomenology.
- **Outlook** (recently developed, under elaboration):
 - EFG is almost completely integrable, can be quantized as almost Kähler–Fedosov/ A–brane geometries, and renormalizable for bi–connection/gauge gravity models.
 - Finsler for black holes (ellipsoids, toruses, holes, wormholes, solitons); anisotropic cosmological models (off–diagonal inflation, dark energy/matter etc).
 - Noncommutative/ Ricci–Finsler flows, emergent (non) commutative Lagrange–Finsler analogous gravity and quantization, Clifford–Finsler algebroids etc.